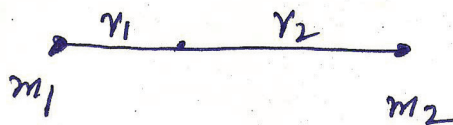


No changeⁱⁿ inter nuclear distance, hence, $V=0$
 There are two angular variables viz., θ & ϕ



$$m_1 r_1 = m_2 r_2$$

$$r_1 + r_2 = R$$

$$KE = \frac{1}{2} m_1 v_1^2 + \frac{1}{2} m_2 v_2^2 = \frac{1}{2} m_1 (r_1 \omega)^2 + \frac{1}{2} m_2 (r_2 \omega)^2$$

$$= \left(\frac{1}{2} m_1 r_1^2 + \frac{1}{2} m_2 r_2^2 \right) \omega^2 = \frac{1}{2} (m_1 r_1^2 + m_2 r_2^2) \omega^2$$

$$= \frac{1}{2} I \omega^2 \quad \text{||| } \frac{1}{2} m v^2 \text{ for linear motion}$$

$$= \frac{L^2}{2I} \quad \text{where; } L = I \omega \quad \text{||| } p = m v$$

$$\hat{H} = \hat{K}E + \hat{P}E = \hat{K}E = \frac{L^2}{2I}$$

L^2 Schrodinger equation $\frac{2I}{\hbar^2} \hat{H} \psi = E \psi \Rightarrow \frac{L^2}{2I} \psi = E \psi$
 in polar coordinates is

$$L^2 = - \frac{\hbar^2}{4\pi^2} \left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right]$$

$$- \frac{\hbar^2}{8\pi^2 I} \left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right] \psi = E \psi$$

$$\text{i.e., } \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2} + \frac{8\pi^2 I}{\hbar^2} E \psi = 0$$

The equation can be solved by the method of separation of variables. The solution expected is of the form

$$\Psi_{\theta, \phi} = \theta(\theta) \cdot \bar{\Phi}(\phi)$$

Substituting for Ψ we get,

$$\frac{\bar{\Phi}}{\sin \theta} \left(\frac{\partial}{\partial \theta} (\sin \theta \frac{\partial \theta}{\partial \theta}) \right) + \frac{\theta}{\sin^2 \theta} \frac{\partial^2 \bar{\Phi}}{\partial \phi^2} + \frac{8\pi^2 I E \theta \bar{\Phi}}{h^2} = 0$$

Multiplying both sides by $\frac{\sin^2 \theta}{\theta \bar{\Phi}}$ we get

$$\frac{\sin \theta}{\theta} \frac{\partial}{\partial \theta} (\sin \theta \frac{\partial \theta}{\partial \theta}) + \frac{1}{\bar{\Phi}} \frac{\partial^2 \bar{\Phi}}{\partial \phi^2} + \frac{8\pi^2 I E \sin^2 \theta}{h^2} = 0$$

$$\frac{\sin \theta}{\theta} \frac{\partial}{\partial \theta} (\sin \theta \frac{\partial \theta}{\partial \theta}) + \frac{8\pi^2 I E \sin^2 \theta}{h^2} = - \frac{1}{\bar{\Phi}} \frac{\partial^2 \bar{\Phi}}{\partial \phi^2}$$

$$= m^2 \text{ (say)}$$

$$\therefore \frac{\partial^2 \bar{\Phi}}{\partial \phi^2} + m^2 \bar{\Phi} = 0$$

----- (1)

$$\frac{\sin \theta}{\theta} \left(\frac{\partial}{\partial \theta} \sin \theta \frac{\partial \theta}{\partial \theta} \right) + \frac{8\pi^2 I E \sin^2 \theta}{h^2} - m^2 = 0$$

----- (2)

The solution for equation-2 is

$$\Phi = e^{\pm im\phi} \quad \text{for } 0 \leq \phi \leq 2\pi$$

$$m = 0, \pm 1, \pm 2, \dots$$

The normalised function Φ is

$$\Phi = \frac{1}{\sqrt{2\pi}} e^{\pm im\phi}$$

Equation-2 can be further modified as follows:

Multiply by $(\frac{\theta}{\sin^2 \theta})$ both sides.

$$\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \theta}{\partial \theta} \right) + \left(\frac{8\pi^2 I E}{h^2} - \frac{m^2}{\sin^2 \theta} \right) \theta = 0$$

$$\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \theta}{\partial \theta} \right) + \left(\beta - \frac{m^2}{\sin^2 \theta} \right) \theta = 0 \quad \dots 3$$

Equation-3 is similar to a Legendre equation "associated" with

$$\beta = l(l+1)$$

\therefore Equation-4 has its solution of the form

$$\theta = P_l^m(x) = (1-x^2)^{m/2} \frac{d^m}{dx^m} P_l(x) \quad \text{Associated Legendre polynomial of degree } l$$

where,

$$l \geq m \quad \& \quad m = 0, 1, 2, 3, \dots$$

$$P_l(x) = \frac{1}{2^l l!} \frac{d^l}{dx^l} (1-x^2)^l, \quad \text{where } x = \cos \theta$$

"Legendre polynomial" of degree l

Moreover

$$\beta = l(l+1) = \frac{8\pi^2 I E}{h^2}$$

$$\therefore E = \frac{h^2}{8\pi^2 I} l(l+1) \quad \text{"Rotational energy" quantised}$$

Hence, $N P_l^m(\cos\theta) \cdot N' e^{\pm im\phi}$ by l .

$$\psi(\theta, \phi) = \Theta \Phi = \left[\frac{2l+1}{2} \frac{(l-m)!}{(l+m)!} \right]^{\frac{1}{2}} P_l^m(\cos\theta) \frac{1}{\sqrt{2\pi}} e^{\pm im\phi}$$

Where N & N' are the normalisation constant as given on allowed values. "Spherical harmonics."

$$l = 0 \quad 1 \quad 1 \quad 2 \quad 2 \quad 2 \quad 3 \quad 3 \quad 3 \quad 3$$

$$m = 0 \quad 0 \quad 1 \quad 0 \quad 1 \quad 2 \quad 0 \quad 1 \quad 2 \quad 3$$

NB: $l \geq |m|$

$$P_0^0(x) = P_0^0(\cos x) = P_0(\cos x) = \frac{1}{\sqrt{2}}$$

$$P_1^0(x) = \left(\frac{3}{2}\right)^{\frac{1}{2}} P_1(\cos x) = \left(\frac{3}{2}\right)^{\frac{1}{2}} \frac{1}{2 \cdot 1!} \frac{d}{dx} (1-x^2)$$

$$= \left(\frac{3}{2}\right)^{\frac{1}{2}} \frac{1}{2} (-2x) = -\left(\frac{3}{2}\right)^{\frac{1}{2}} \cos x$$